

Global F-Theory Compactifications with Higher Rank Abelian Symmetries

Mirjam Cvetič



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arXiv:1303.6970 [hep-th]: M. C. Denis Klevers, Hernan Piragua

arXiv:1306.0236 [hep-th]: M. C. A. Grassi, D. Klevers, H. Piragua

arXiv:130n.nnnn [hep-th] (UPR-1251-T): M.C., D. Klevers, H. Piragua

Also: Xiv:1210.6034 [hep-th]: M. C., Thomas W. Grimm, D. Klevers



F-theory Compactifications with additional $U(1)$'s

MOTIVATION

Why F-theory Compactification?

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Domain of string theory landscape with promising particle physics

- Focus D=4 N=1 SUSY GUT's [SU(5), SO(10)]
w/chiral matter, Yukawa couplings 10 10 5,...

➡ GUT-model building in F-theory

- **Moduli stabilization** (fluxes) [Gukov, Vafa, Witten],...

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Local: [Donagi, Wijnholt; Beasley, Heckman, Vafa;
... Review: Heckman, ...]

Global: [Marsano, Saulina, Schäfer-Nameki;
Blumenhagen, Grimm, Jurke, Weigand; ...
M.C., Halverson, Garcia-Etxebarria; ...]

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M.C., Halverson, Garcia-Etxebarria; ...]


Conceptual: geometric description at finite string coupling

[Vafa; Witten; ...]

- F-theory via finite coupling Type IIB string theory:
Consistent set-up of back-reacted seven-branes
Non-perturbative coupling regions on non-Calabi-Yau geometry
- F-theory via Geometry:
Globally defined elliptically fibered Calabi-Yau manifold

Why Abelian Symmetries in F-theory?

Particle physics: important ingredient of Beyond Standard Model Physics

- Light $U(1)$ gauge bosons: Z' -physics, NMSSM, $U(1)_{PQ}$, ...
- Massive (Stückelberg) $U(1)$ gauge bosons: low energy global symmetry
  selection rules (proton decay; R-parity violation; neutrino masses...)

Multiple $U(1)$'s desirable

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Conceptual: new types of elliptic fibrations

- Related to Abelian **Mordell-Weil group** of elliptic fibrations
➡ torsion part studied Torsion part: [Morrison, Vafa;
Aspinwall, Morrison;...]
w/free part less understood (global issues) For toric K3: [Grassi, Perduca]
- Few systematic studies in contrast to non-Abelian groups
Non-Abelian: [Kodaira; Tate;
Morrison, Vafa; Bershadsky et al.;...]

Outline & Summary of the talk

- I. Construction of elliptically fibered Calabi-Yau manifolds w/
rank 2 Mordell-Weil (MW) group
- II. Determination of matter representations and multiplicity in D=6 and D=4
- III. First construction of G_4 -fluxes on Calabi-Yau four-folds with rk=2 MW-group
- IV. Construction of explicit $U(1) \times U(1)$ and $SU(5) \times U(1) \times U(1)$ w/spectra & chirality

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Two-fold advances: Geometry & M-theory/F-theory duality

- I. Geometry: in D=6 determines all matter representations and multiplicity
in D=4 G_4 fluxes & some of matter surfaces identified \rightarrow some chiralities
- II. M-theory/F-theory duality in D=3: constraints on $G_4 \rightarrow$ Chern-Simons terms determine
 - all chiral indices (tested against geom. calc.)
 - confirm cancellation of all anomalies

The Type IIB perspective

F-THEORY HIGHLIGHTS

The Type IIB perspective

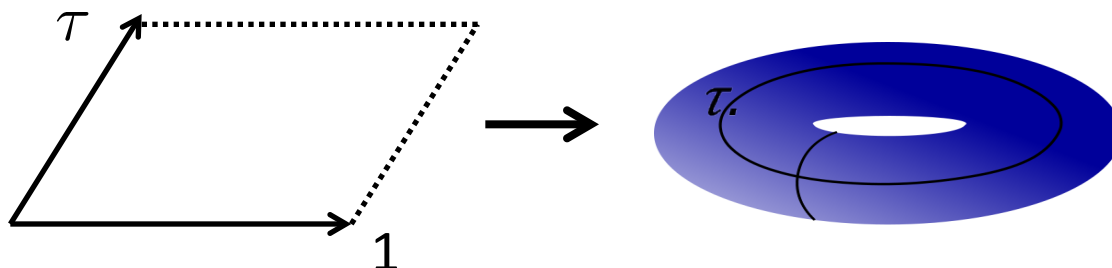
F-THEORY HIGHLIGHTS

[At Strings'12: D-instantons in F-theory; Heterotic M-theory perspective]

[M.C., Donagi, Halverson, Marsano]

F-theory via Type IIB: basic ingredients

- F-theory is a geometric, $SL(2, \mathbb{Z})$ invariant formulation of Type IIB string theory: invariant geometric object is two-torus $T^2(\tau)$ [Vafa]



- Modular parameter τ of $T^2(\tau)$: $\tau \equiv C_0 + ig_s^{-1}$ Type IIB axion-dilaton ($SL(2, \mathbb{Z}) = S$ -duality)

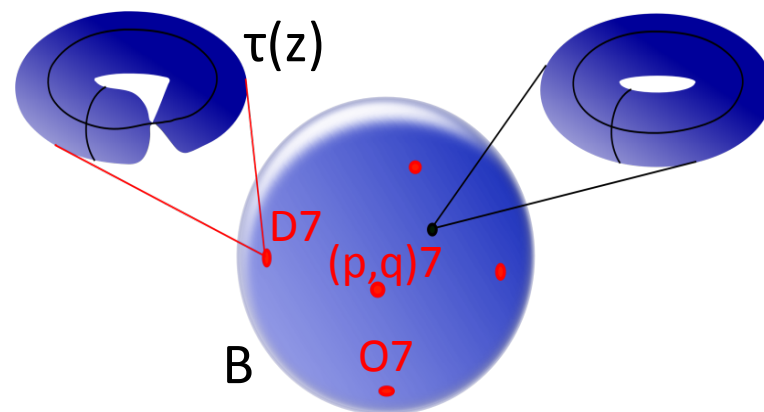
- $T^2(\tau)$ -fibration over a base space B :

Weierstrass parameterization:

$$y^2 = x^3 + fxz^4 + gz^6$$

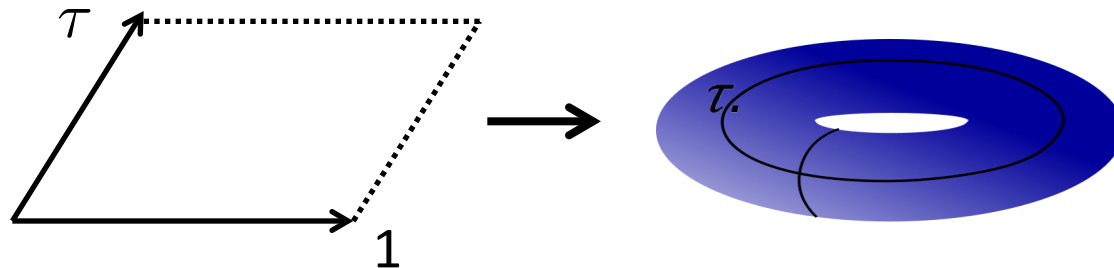
f, g - function fields on B

$[z:x:y]$ homog. coords on $\mathbf{P}^2(1,2,3)$



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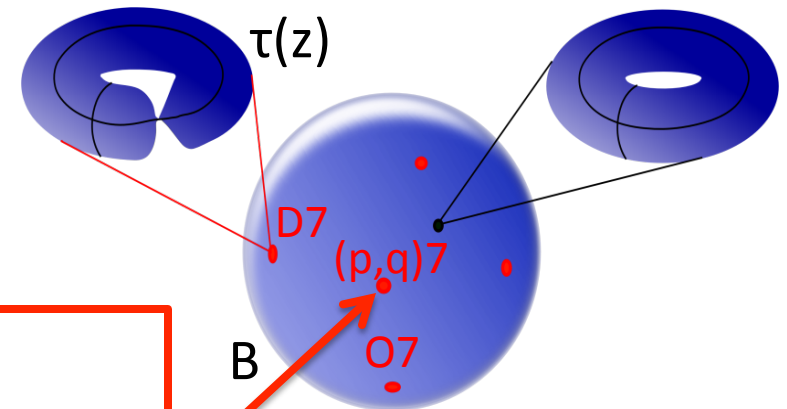


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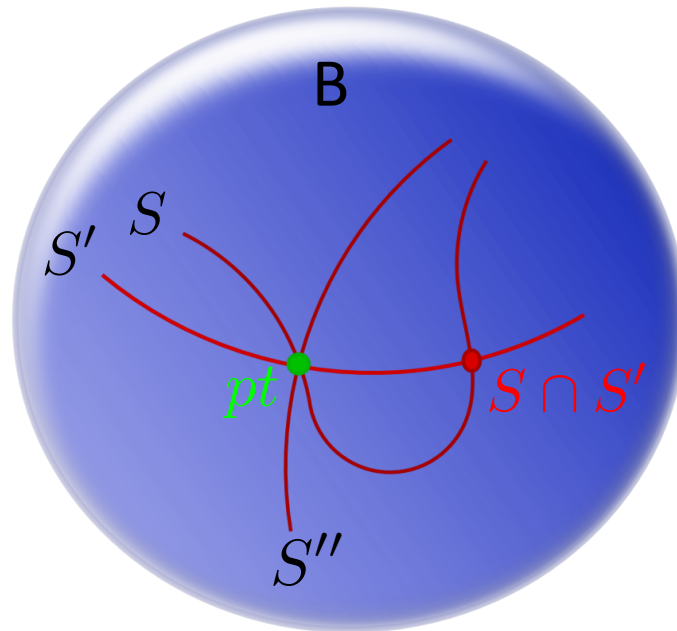
7-branes

non-perturbative regime:

$$g_s \rightarrow \infty \longleftrightarrow \text{singular } T^2(\tau)$$

F-theory via Type IIB: basic ingredients

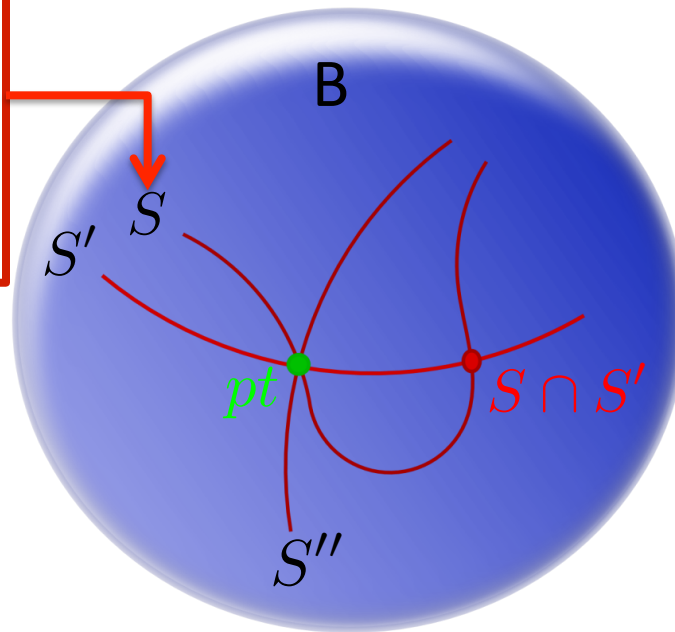
- Total space of $T^2(\tau)$ -fibration: singular elliptic Calabi-Yau manifold X
D=4, N=1 vacua: fourfold X_4 [D=6, N=1 vacua: threefold X_3]
- X-singularities encode complicated set-up of intersecting 7-branes:



F-theory via Type IIB: basic ingredients

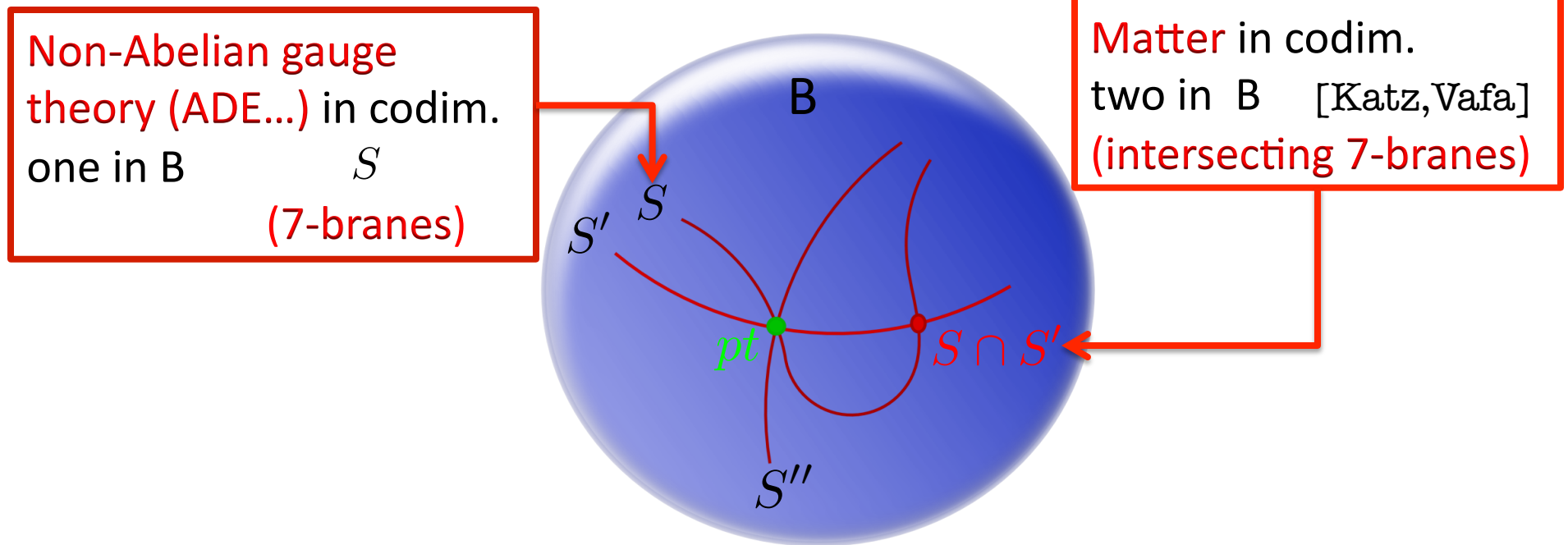
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Non-Abelian gauge
theory (ADE...) in codim.
one in B
 S
(7-branes)



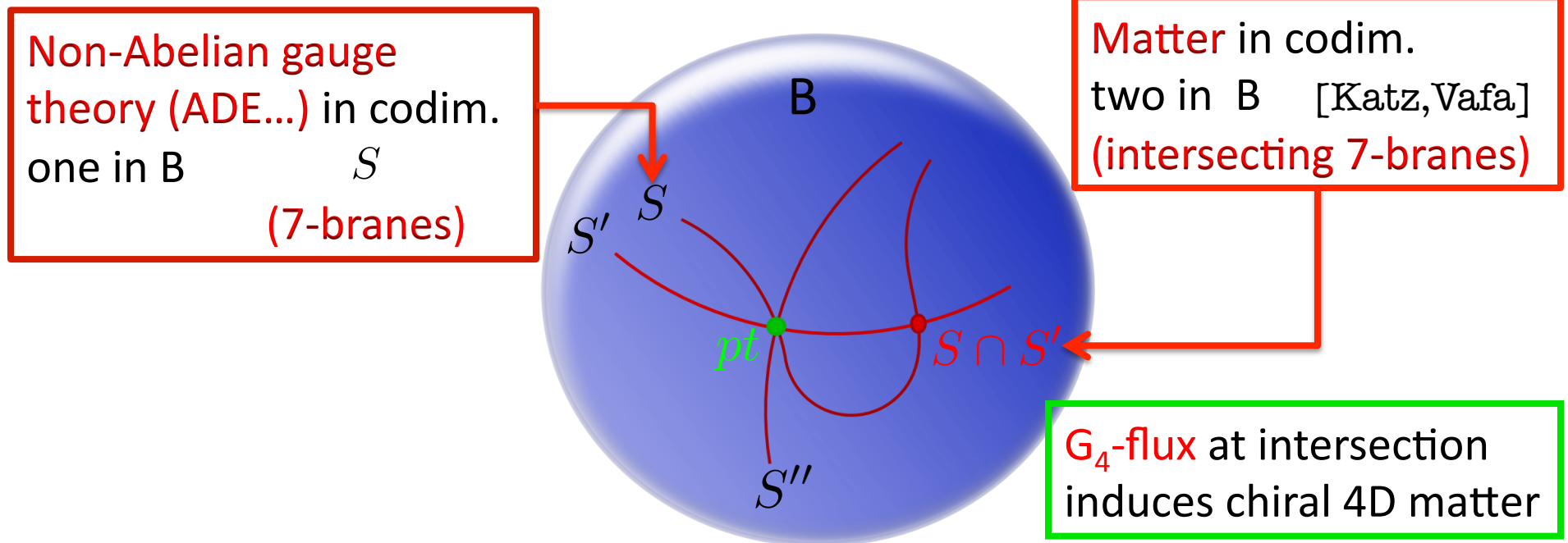
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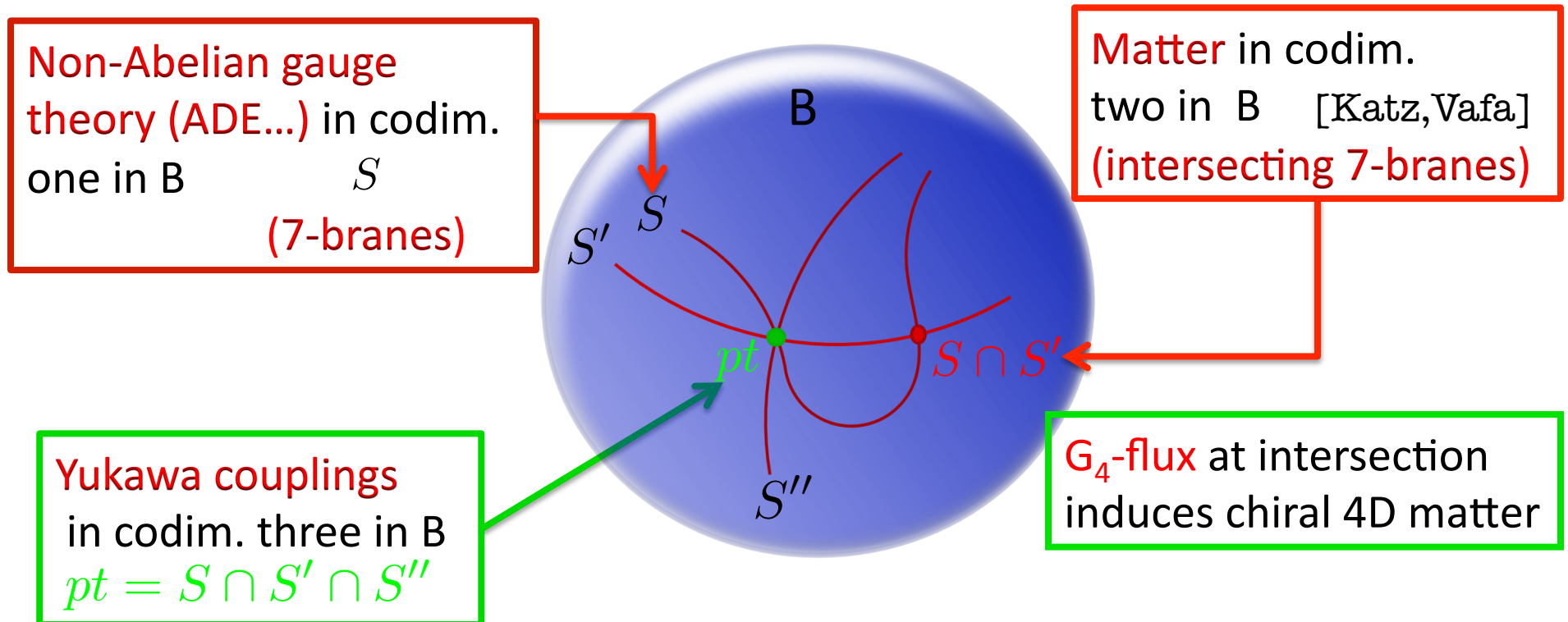
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F-theory via Type IIB: basic ingredients

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 $D=4, N=1$ vacua: fourfold X $[D=6, N=1$ vacua: threefold $X]$
- X -singularities encode complicated set-up of intersecting 7-branes:



Constructing elliptic fibrations with rank two Mordell-Weil groups

U(1)XU(1) SYMMETRY IN F-THEORY

MW-group of rational sections & U(1)'s

4D Abelian gauge fields arise from classical Kaluza-Klein-reduction of C_3

$$C_3 = A^B \omega_B \supset A^i \omega_i + A^m \omega_m$$

(1,1)-forms on X Cartans of non-Abelian group U(1)-gauge fields

The diagram illustrates the decomposition of the 3-form field C_3 into two parts. The first part, $A^B \omega_B$, is identified as (1,1)-forms on X. The second part, $A^i \omega_i + A^m \omega_m$, is identified as the Cartans of a non-abelian group and U(1)-gauge fields. Red arrows point from the text labels to the corresponding terms in the equation.

MW-group of rational sections & U(1)'s

4D Abelian gauge fields arise from **classical Kaluza-Klein-reduction of C_3**

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(1,1)-forms on X
Cartans of non-Abelian group
U(1)-gauge fields

Construction of (1,1)-form ω_m via rational sections

1. **Rational point** Q on elliptic curve E with zero point P

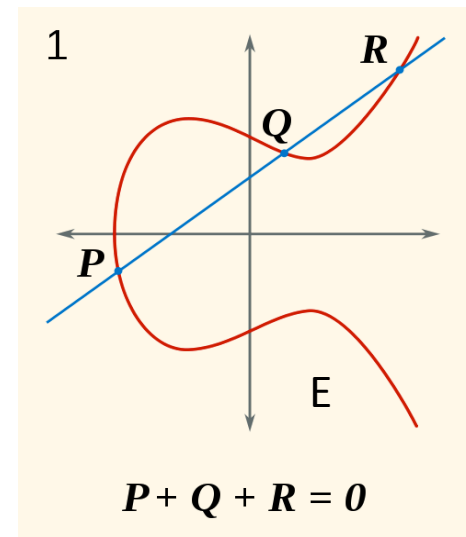
- is solution $[z_Q : x_Q : y_Q]$ in field K of Weierstrass form

$$y^2 = x^3 + fxz^4 + gz^6$$

- Rational points **form group** (addition) on E



Mordell-Weil group of rational points



[wikipedia.org]

2. Q induces **rational section** $\hat{s}_Q : B \rightarrow X$ of the fibration

(1,1)-form ω_m Poincaré dual to divisor class S_Q (related to \hat{s}_Q via Shioda map)

Construction of elliptic curve with $\text{rk}(\text{MW})=2$

[M.C., Klevers, Piragua]

Elliptic curve E with two rational points Q, R

related work: [Borchman, Mayrhofer, Weigand]

$\text{rk}(\text{MW})=1$: [Morrison, Park; Mayrhofer, Palti, Weigand]

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Elliptic curve E with two rational points Q, R

Consider line bundle $M=O(P+Q+R)$ of degree 3 on E (non-generic cubic in \mathbf{P}^2)

➔ natural representation as hypersurface $p=0$ in del Pezzo dP_2

$$p = u(s_1 u^2 e_1^2 e_2^2 + s_2 u v e_1 e_2^2 + s_3 v^2 e_2^2 + s_5 u w e_1^2 e_2 + s_6 v w e_1 e_2 + s_8 w^2 e_1^2) + s_7 v^2 w e_2 + s_9 v w^2 e_1$$

$[u:v:w:e_1:e_2]$ –homogeneous coordinates of dP_2

(blow-up of \mathbf{P}^2 w/ $[u':v':w']$ at 2 points: $u'=ue_1e_2, v'=ve_2, w'=we_1$)

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u v w e₁ e₂

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Points represented by intersections of different divisors in dP_2 with p

Classification of dP_2 elliptic fibrations

[M.C., Klevers, Piragua; M.C., Grassi, Klevers, Piragua]

I. Ambient space:

- dP_2 fibration determined by two divisors \mathcal{S}_7 and \mathcal{S}_9 (loci of $s_7=0, s_9=0$)

$$\begin{array}{ccc} dP_2 & \longrightarrow & dP_2^B(\mathcal{S}_7, \mathcal{S}_9) \\ & & \downarrow \\ & & B \end{array}$$

II. Calabi-Yau hypersurface X:

- cuts out E in dP_2
- coefficients s_i in CY-equation get lifted to sections of the base B (only s_7, s_9 independent)
- coordinates $[u:v:w:e_1:e_2]$ lifted to sections

$$\begin{array}{ccc} \hat{E} \subset dP_2 & \longrightarrow & X \\ & \nearrow \text{sections } \hat{s}_P, \hat{s}_Q, \hat{s}_R & \downarrow \\ & & B \end{array}$$

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Birational map to Weierstrass fibration explicitly worked out

Classification of dP_2 elliptic fibrations

Construction of general elliptic fibrations:

section	bundle	section	bundle
u'	$\mathcal{O}(H - E_1 - E_2 + \mathcal{S}_9 + [K_B])$	s_1	$\mathcal{O}(3[K_B^{-1}] - \mathcal{S}_7 - \mathcal{S}_9)$
v'	$\mathcal{O}(H - E_2 + \mathcal{S}_9 - \mathcal{S}_7)$	s_2	$\mathcal{O}(2[K_B^{-1}] - \mathcal{S}_9)$
w'	$\mathcal{O}(H - E_1)$	s_3	$\mathcal{O}([K_B^{-1}] + \mathcal{S}_7 - \mathcal{S}_9)$
e_1	$\mathcal{O}(E_1)$	s_5	$\mathcal{O}([2K_B^{-1}] - \mathcal{S}_7)$
e_2	$\mathcal{O}(E_2)$	s_6	$\mathcal{O}([K_B^{-1}])$
— CY-condition: \mathcal{S}_7 and \mathcal{S}_9 fixed		s_7	$\mathcal{O}(\mathcal{S}_7)$
		s_8	$\mathcal{O}([K_B^{-1}] + \mathcal{S}_9 - \mathcal{S}_7)$
		s_9	$\mathcal{O}(\mathcal{S}_9)$

Engineer non-Abelian groups: make s_i non-generic

Can apply to toric cases w/ two $U(1)$'s

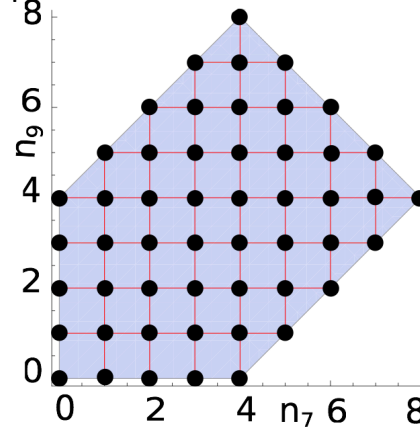
[Bonetti, Braun, Grimm, Hohenegger; Borchmann, Mayrhofer, Palti, Weigand;
Braun, Grimm, Keitel]

Classification of dP_2 elliptic fibrations

All topologically distinct $D=6$ & $D=4$ vacua for fixed base B . Example:

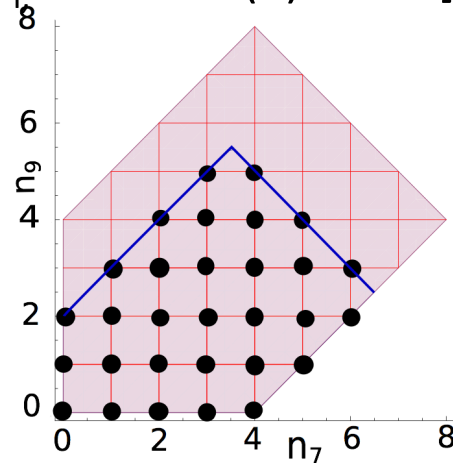
1. $B=\mathbf{P}^3$, X generic [all s_i exist, generic]: $U(1) \times U(1)$

$$\begin{aligned} S_7 &= n_7 H_{\mathbb{P}^3} \\ S_9 &= n_9 H_{\mathbb{P}^3} \end{aligned}$$



2. $B=\mathbf{P}^3$, X non-generic [s_i realize $SU(5)$ at $t=0$]: $SU(5) \times U(1) \times U(1)$

$$\begin{aligned} s_1 &= t^3 s'_1 \\ s_2 &= t^2 s'_2 \\ s_3 &= t^2 s'_3 \\ s_5 &= t s'_5 \end{aligned}$$



[M.C., Klevers, Piragua]

Can construct and **study all these CYs explicitly**
(no restriction to **toric** hypersurfaces seems necessary)

Codimension two singularities of dP_2 -elliptic fibrations

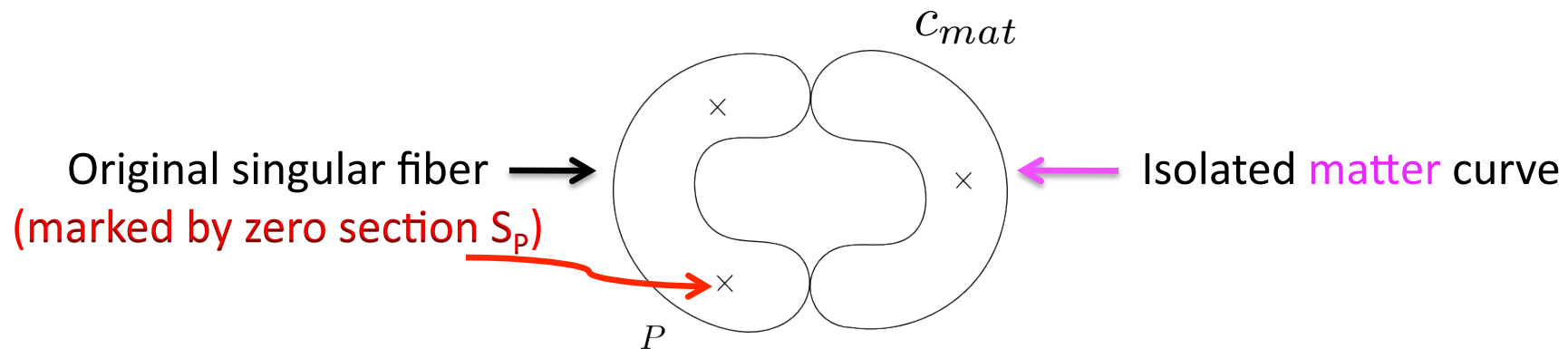
MATTER $U(1) \times U(1)$ F-THEORY VACUA

Matter representations

[M.C., Klevers, Piragua]

related work: [Borchman, Mayrhofer, Weigand]

- Matter in F-theory arises from a co-dimension two singularities in B
- Singular fiber resolved into reducible curves $E = c_1 + c_{\text{mat}}$ w/ $c_1 \cdot c_{\text{mat}} = 2$
(c_{mat} - M2-branes wrapping isolated \mathbb{P}^1 in reducible fiber)



Advances in higher co-dimension singularities: [Esole, Yau].., [Lawrie, Schäfer-Nameki]

Recent advances via deformations of singularities: [Halverson, Grassi, Shaneson]

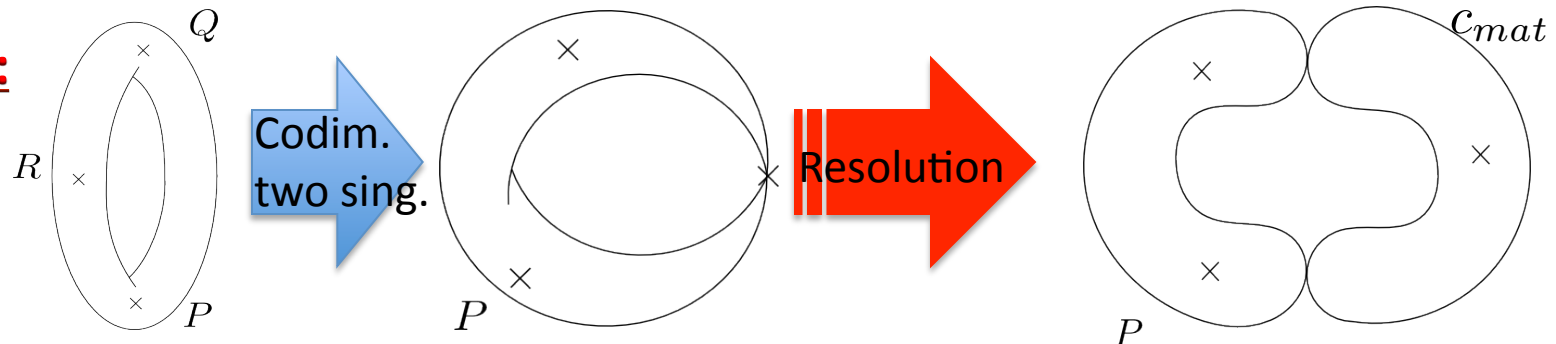
Charged matter of Type I

Charge formula:

$$q_1 = (S_Q - S_P) \cdot c_{mat} \quad q_2 = (S_R - S_P) \cdot c_{mat}$$

Strategy: look for collisions of rational sections with singularities in Weierstrass fibration (WSF)

WSF:

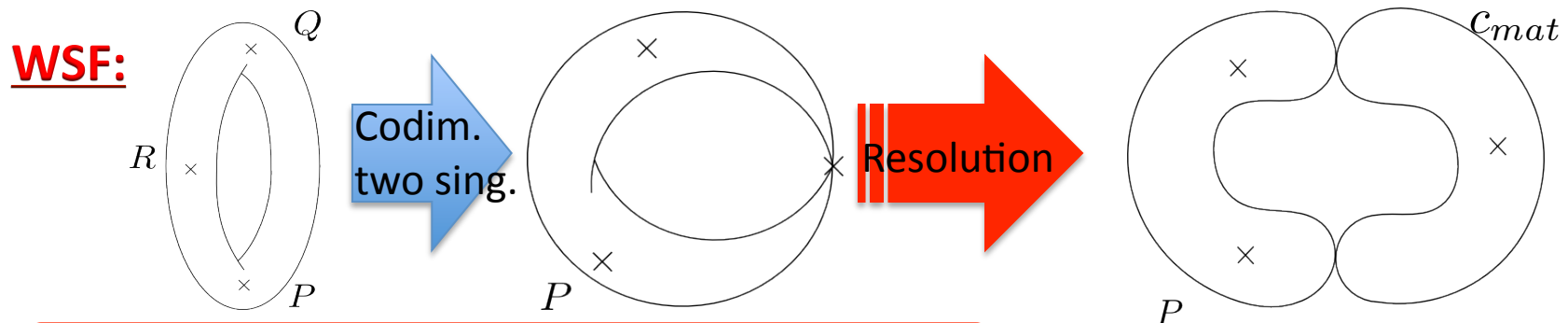


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List of charged matter representations

Representation

Collision pattern

i. $(q_1, q_2) = (1, 0)$

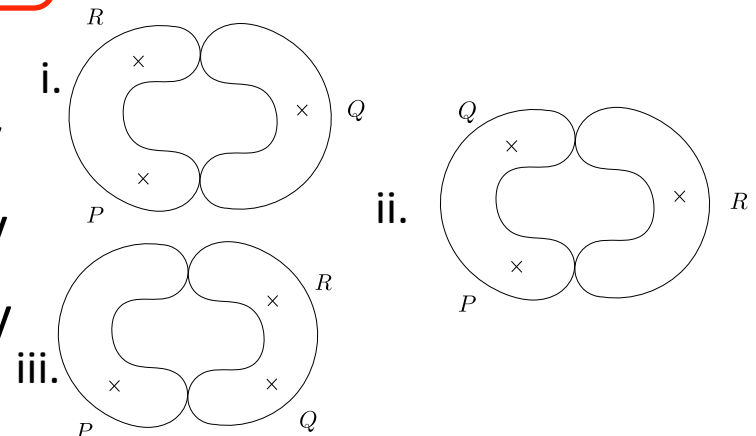
$Q \rightarrow$ WS-singularity

ii. $(q_1, q_2) = (0, 1)$

$R \rightarrow$ WS-singularity

iii. $(q_1, q_2) = (1, 1)$

$Q, R \rightarrow$ WS-singularity



Charged matter of Type II

Strategy: look for loci in B where the sections are ill-defined

$$\begin{aligned} P : E_2 \cap p &= [-s_9 : s_8 : 1 : 1 : 0], & Q : E_1 \cap p &= [-s_7 : 1 : s_3 : 0 : 1], \\ R : D_u \cap p &= [0 : 1 : 1 : -s_7 : s_9]. \end{aligned}$$

sections no longer points in E, wrap entire \mathbf{P}^1 in smooth X

Charged matter of Type II

Strategy: look for loci in B where the **sections are ill-defined**

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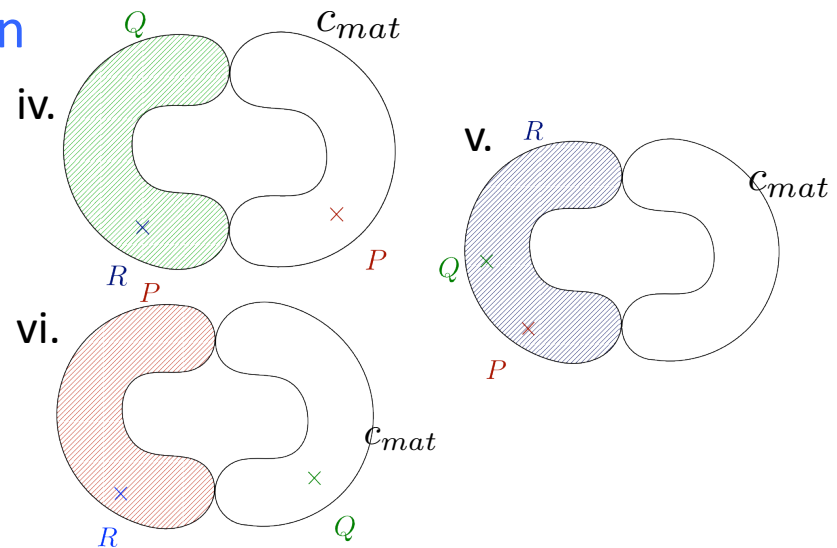
List of charged matter representations

Representation	Ill-defined section
----------------	---------------------

iv. $(q_1, q_2) = (-1, 1)$	Q at $s_3 = s_7 = 0$
----------------------------	----------------------

v. $(q_1, q_2) = (0, -2)$	R at $s_7 = s_9 = 0$
---------------------------	----------------------

vi. $(q_1, q_2) = (-1, -2)$	zero section P at $s_8 = s_9 = 0$
-----------------------------	--------------------------------------



Summary of Matter Representations

	$U(1) \times U(1)$
Type I	$(1, 0) \ (0, 1) \ (1, -1)$
Type II	$(-1, 1) \ (0, 2) \ (-1, -2)$

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Apply analogous techniques to determine matter representations

Specific example:

$$\begin{aligned} s_1 &= t^3 s'_1 \\ s_2 &= t^2 s'_2 \\ s_3 &= t^2 s'_3 \\ s_5 &= t s'_5 \end{aligned}$$

w/ $SU(5)$ at $t=0$

Summary of Matter Representations

	$U(1) \times U(1)$	$SU(5) \times U(1) \times U(1)$
Type I	$(1, 0) (0, 1) (1, -1)$	$(\mathbf{5}, -\frac{2}{5}, 0) (\mathbf{5}, \frac{3}{5}, 0) (\mathbf{5}, -\frac{2}{5}, -1)$
Type II	$(-1, 1) (0, 2) (-1, -2)$	$(\mathbf{5}, -\frac{2}{5}, 1) (\mathbf{5}, \frac{3}{5}, 1) (\overline{\mathbf{10}}, -\frac{1}{5}, 0)$



X non-generic \rightarrow realize $SU(5) \times U(1) \times U(1)$

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Specific example:

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w/ $SU(5)$ at $t=0$

Matter multiplicities

MATTER SPECTRUM IN 6D

6D matter multiplicities

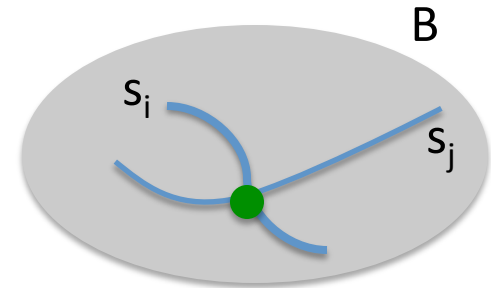
[M.C.,Klevers,Piragua]

Matter multiplicities = number of points in codimension 2 matter loci in B

1. Matter of Type II: simple complete intersection

$$s_i = s_j = 0$$

Number of points = $\deg(s_i) * \deg(s_j)$



6D matter multiplicities

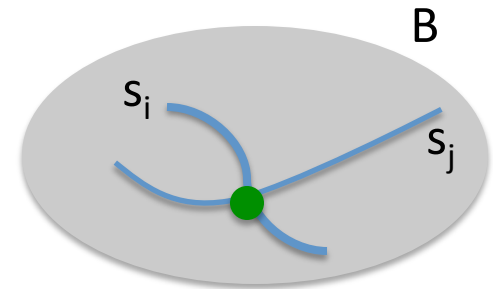
[M.C., Klevers, Piragua]

Matter multiplicities = number of points in codimension 2 matter loci in B

1. Matter of Type II: simple complete intersection

$$s_i = s_j = 0$$

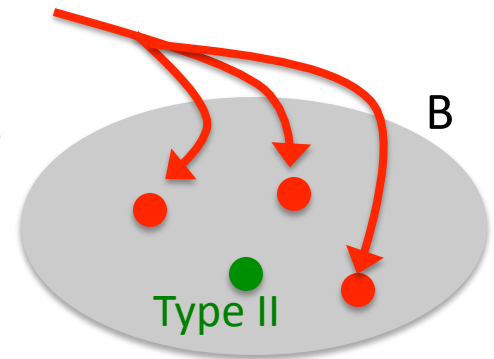
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2. Matter of Type I: opposite of complete intersections

Described by prime ideals (8 polynomial equations)

Counting of points via resultant of polynomial system



6D matter multiplicities

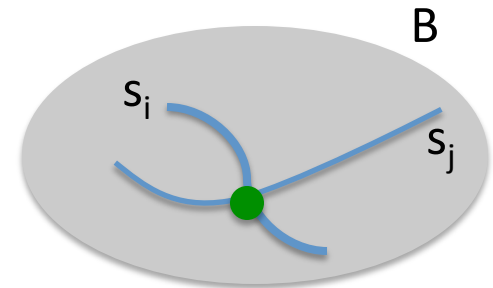
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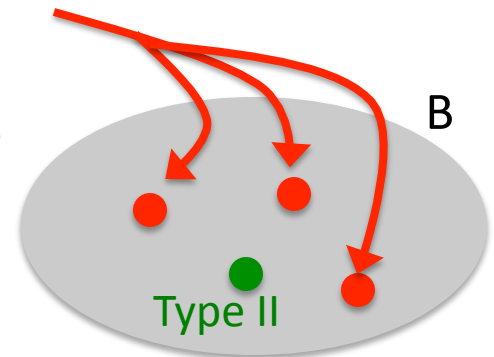
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2. Matter of Type I: opposite of complete intersections

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Method general! Can now apply to examples of

rk=1 MW [Morrison, Park; Mayrhofer, Palti, Weigand]

rk=2 MW [Borchman, Mayrhofer, Weigand], ...

6D matter spectrum & multiplicities

6D matter spectrum and multiplicities can be obtained over any base B

Example: $B = \mathbf{P}^2$ $w/U(1) \times U(1)$

	(q_1, q_2)	Multiplicity
Type I	$(1, 0)$	$54 - 15n_9 + n_9^2 + (12 + n_9)n_7 - 2n_7^2$
	$(0, 1)$	$54 + 2(6n_9 - n_9^2 + 6n_7 - n_7^2)$
	$(1, 1)$	$54 + 12n_9 - 2n_9^2 + (n_9 - 15)n_7 + n_7^2$
Type II	$(-1, 1)$	$n_7(3 - n_9 + n_7)$
	$(0, 2)$	n_9n_7
	$(-1, -2)$	$n_9(3 + n_9 - n_7)$

Integers n_7, n_9 specify all dP_2 -fibration over \mathbf{P}^2

Full spectrum and multiplicities also with $SU(5) \times U(1) \times (1)$ group

Consistency check: spectrum found to cancel 6D anomalies!

Matter surfaces, G_4 -flux & 3D CS-terms

MATTER SPECTRUM IN 4D

4D matter spectrum

[M.C.,Grassi,Klevers,Piragua]

4D-matter representations the same as in 6D (all in the fiber)

4D matter chiralities = codimension two matter loci in B + G_4 -flux:

$$\chi(\mathbf{R}) = -\frac{1}{4} \int_{\mathcal{C}_{\mathbf{R}}} G_4$$

Geometry: I. Matter surfaces:

points in $B_2 \rightarrow$ matter curves $\Sigma_{\mathbf{R}}$ in B_3

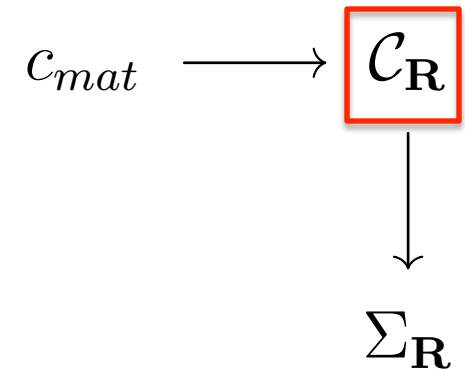
Type II matter surfaces found

Type I matter-hard

II. G_4 -flux:

Construction of homology $H_V^{(2,2)}(\hat{X})$

First construction of G_4 -flux with non-holomorphic zero-section



4D matter spectrum

[M.C., Grassi, Klevers, Piragua]

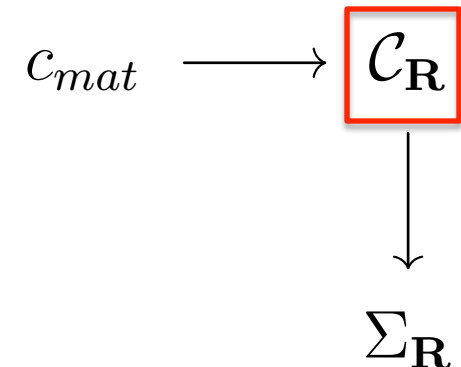
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Type II matter surfaces found

Evaluate integrals

Type I matter-hard

Chiral index

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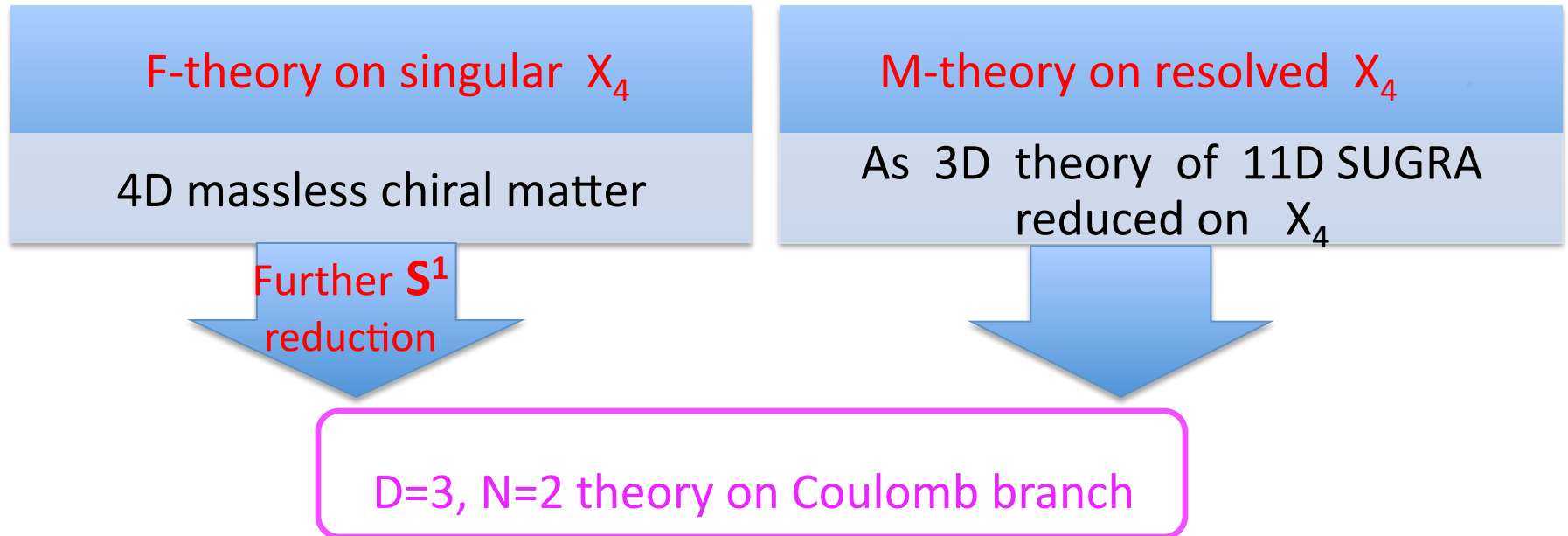
Conditions for G_4 -flux in F-theory

G_4 -flux in F-theory = G_4 -flux in M-theory + extra conditions



M/F-theory duality in D=3

c.f., John Schwarz's talk



c.f., Nati Seiberg's talk

Match as quantum effective actions in IR
(Integrate out massive states: massive 4D matter, KK-states)

Conditions for G_4 -flux in F-theory

I. G_4 in M-theory: 3D Cherns-Simons terms are classical

$$S_{CS}^{3D} = \int \frac{1}{2} \Theta_{AB} A^A \wedge F^B \quad \Theta_{AB} = \int_{\hat{X}_4} G_4 \wedge \omega_A \wedge \omega_B$$

D_A = basis of divisors on X_4

II. G_4 in F-theory (3D Coulomb branch):

some classical: 4D gaugings of RR-axions (GS) [Grimm,Kerstan,Palti,Weigand]

some exotic (set to zero) [Grimm,Savelli]

some loop-generated: massive fermions on 3D Coulomb branch + KK-states

$$\Theta_{AB}^{\text{loop}} = \frac{1}{2} \sum_{\mathbf{R}} \chi(\mathbf{R}) \sum_{q \in \mathbf{R}} \sum_k q_A q_B \text{sign}(m_{CB} + \frac{k}{r_{\text{KK}}})$$

[Aharony,Hanany,Intriligator,Seiberg,Strassler]

...

Conditions for G_4 -flux in F-theory

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G_4 -conditions

Constrain G_4 in M-theory (I.) so that $\Theta_{AB}=0$ for CS-terms that in F-theory (II.) are neither classically, nor one-loop generated 

Nonzero Θ_{AB} in turn determine all chiralities and all anomaly cancellations!
...[Grimm,Hayashi;M.C.,Grassi,Klevers,Piragua] [M.C.,Grimm,Klevers]

The full 4D spectrum

Example $B=P^3$ w/ $U(1) \times U(1)$: most general solution for G_4 -flux

$$G_4 = a_5 n_9 (4 - n_7 + n_9) H_B^2 + 4a_5 H_B S_P + a_3 H_B \sigma(\hat{s}_Q) + a_4 H_B \sigma(\hat{s}_R) + a_5 S_P^2$$

(q_1, q_2)	4D chiralities
$(1, 0)$	$\frac{1}{4} [a_5 n_7 n_9 (4 - n_7 + n_9) + a_3 (2n_7^2 - (12 - n_9)(8 - n_9) - n_7(16 + n_9))]$
$(0, 1)$	$\frac{1}{2} [a_5 n_9 (4 - n_7 + n_9) (12 - n_9) - a_4 (n_7 (8 - n_7) + (12 - n_9)(4 + n_9))]$
$(1, 1)$	$\frac{1}{4} [2a_5 n_9 (4 - n_7 + n_9) (12 - n_9) - (a_3 + a_4) (n_7^2 + n_7 (n_9 - 20) + 2(12 - n_9)(4 + n_9))]$
$(-1, 1)$	$\frac{1}{4} (a_3 - a_4) n_7 (4 + n_7 - n_9)$
$(0, 2)$	$\frac{1}{4} n_7 n_9 (-2a_4 + a_5 (4 - n_7 + n_9))$
$(-1, -2)$	$-\frac{1}{4} n_9 (n_7 - n_9 - 4) (a_3 + 2a_4 + a_5 (n_7 - 2n_9))$

All 4D anomalies cancelled;

Chiralities checked against Type II matter geometric chirality calculations

The full 4D spectrum

Example $B=P^3$ w/ $U(1) \times U(1)$: most general solution for G_4 -flux

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Same methods for $SU(5) \times U(1) \times U(1)$ applied:

G_4 -flux has 7 parameter; all 4D chiralities determined; anomalies checked;

Chirality checked against Type II matter geometric calculations



Summary

- Systematic construction of elliptic fibrations with $\text{rk}=2$ MW-groups
- $D=6$: Matter spectrum and multiplicity for general B
 $U(1) \times U(1) \times SU(5) \times U(1) \times U(1)$ - All Geometry
- $D=4$ Matter spectrum and chirality
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(new CS-terms from charged KK-states)
Determine all chiralities (checked against geom. calc.)

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Outlook

- 4D: Generalization to other bases, SUSY, ...  Phenomenology
- More $U(1)$'s.. 

Announce: Elliptic CY with $\text{rk}(\text{MW})=3$

[M.C., Klevers, Piragua, Peng Song] to appear

Elliptic curve E with three rational points Q, R, S

Line bundle $M=\mathcal{O}(P+Q+R+S)$ of degree 4 on E (non-generic biquadric in \mathbf{P}^3)

➡ Generic E : Calabi-Yau Complete Intersection (defined by two equations) in the blow-up of \mathbf{P}^3 at three points

-The birational map to the Weierstrass model worked out

-Elliptic fibration, classification

-Matter, Multiplicities...

} Work in progress